

Published on Web 03/11/2004

## The Rotational Spectrum and Geometrical Structure of Thiozone, S<sub>3</sub>

Michael C. McCarthy,\* Sven Thorwirth, Carl A. Gottlieb, and Patrick Thaddeus

Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, Massachusetts 02138, and Division of Engineering and Applied Sciences, Harvard University, Cambridge, Massachusetts 02138

Received January 20, 2004; E-mail: mccarthy@cfa.harvard.edu

Elemental sulfur has the largest number of allotropes of any element in the periodic table, and has been the subject of many experimental and theoretical studies. Small sulfur clusters  $(S_2-S_8)$ serve as important benchmarks for the accuracy of quantum chemical calculations because their structures and relative stabilities are sensitive to effects such as electron correlation.<sup>1</sup> Such clusters are also of astronomical interest: they are thought to play important roles in the volcanically driven atmosphere of Io,<sup>2</sup> the innermost Galilean moon of Jupiter, in the photochemistry of the lower atmosphere of Venus,<sup>3</sup> and in the circumstellar shells of carbonrich evolved stars.<sup>4</sup>

Our knowledge of the electronic ground state and geometrical structure of small sulfur clusters is fragmentary. It is well known that S<sub>2</sub>, like O<sub>2</sub>, has a triplet ground state  $(^{3}\Sigma^{-})$ , but there is little direct experimental information on larger clusters from S<sub>3</sub> to S<sub>5</sub>. In contrast, the structure of ozone O<sub>3</sub> has been known for more than 50 years.<sup>5</sup> Larger sulfur clusters  $S_n$ , n > 5, unlike oxygen clusters of the same size, occur as stable monocyclic rings, and many of these (n = 6-8, 10-13, etc.) have been well-characterized by X-ray crystallography and infrared spectroscopy.<sup>6</sup> Several sulfur clusters, starting at S<sub>3</sub>, are calculated to be polar, with dipole moments of a few tenths of one Debye, and therefore are good candidates for detection by Fourier transform microwave (FTM) spectroscopy, where rotational line intensities are proportional to the first power of the dipole moment, not its square, as in conventional spectroscopy. With this technique, we have now detected S<sub>3</sub> and by isotopic substitution determined its geometrical structure to high precision.

Thiozone is stable at fairly high temperature, and its optical and vibrational spectra have been detected in solid argon<sup>7</sup> and in the gas phase<sup>8</sup> at low-resolution, but high-resolution observations to determine its structure have not been reported.<sup>9</sup> Ab initio calculations indicate that the most stable isomer is bent with the  $C_{2v}$  symmetry shown in Figure 1 and that the electronic ground state is singlet <sup>1</sup>A<sub>1</sub>. A nonpolar singlet ring isomer with  $D_{3h}$  symmetry is calculated to lie only 5–10 kcal/mol higher in energy,<sup>1,10,11</sup> and the lowest triplet state, <sup>3</sup>A<sub>1</sub> with  $C_{2v}$  symmetry, is calculated to lie at 14–18 kcal/mol above ground. The calculated dipole moment of the singlet isomer in Figure 1 is 0.51 D.<sup>12</sup>

The structure of thiozone has proven to be a computational challenge because the S–S bond length is quite sensitive to the level of theory, the size of the basis set employed, and the participation of the low-lying 3d orbitals in bonding. More than 30 theoretical studies have been published, ranging from Hartree–Fock to high-level coupled cluster calculations [CCSD(T)]. The bond, for example, is calculated to be as short as 1.888 Å (HF/3-21G\*) or as long as 1.980 Å (BLYP/DZP).<sup>13</sup> The apex angle is generally less sensitive to the level of calculation, with different techniques yielding values in the range 117–118°.

We detected the rotational spectrum of  $S_3$  by means of FTM spectroscopy of a molecular beam — the same technique recently



**Figure 1.** Experimental structure of thiozone, an asymmetric top molecule with  $C_{2\nu}$  symmetry. Estimated uncertainties in units of the last significant digit are given in parentheses. The bond order is only approximate.

used to identify the structurally similar silicon trimer<sup>14</sup> and more than 110 other reactive molecules,<sup>15</sup> including a large number of organic radicals and carbenes. Reactive molecules are created in the throat of a small supersonic nozzle by applying a low-current dc discharge to a short gas pulse created by a fast mechanical valve, the gas in the present work being either hydrogen sulfide H<sub>2</sub>S (1.0%) heavily diluted in a neon or argon buffer, or sulfur vapor (0.1%) in neon which was produced by flowing the buffer gas over sulfur powder resistively heated to 185 °C.

A search for  $S_3$  was undertaken in the frequency band 20–25 GHz, where several strong rotational lines are expected. A fairly large search in frequency was required, because S<sub>3</sub> is not strongly bent, and its moment of inertia along the a-axis is quite sensitive to the molecular structure; even small changes in the bond length (i.e.,  $\pm 0.02$  Å) or apex angle ( $\pm 0.2^{\circ}$ ) can shift rotational lines by 10-20% in frequency. Several tests quickly established that the carrier of one of the lines near 22 GHz only contains sulfur. Because of the high signal-to-noise, the corresponding lines  $(2_{1,1}-2_{0,2})$  of the two mono-34S isotopic species were then readily detected, despite the low fractional abundance of this isotope (4.2%). The close agreement (to 1%) between the observed frequency shifts of the isotopic species and those predicted from the geometry in Figure 1 is strong evidence that the assigned lines are produced by S<sub>3</sub> and no other sulfur cluster. Although H2S was used at first as a precursor gas because it is a good source for sulfur chains such as  $HS_2$ ,  $H_2S_2$ , and H<sub>2</sub>S<sub>3</sub><sup>16</sup> it was then found that sulfur vapor in a neon buffer yielded 25 times stronger lines of S<sub>3</sub>.

A total of 17 *b*-type transitions of  ${}^{32}S_3$  up to J = 12 and  $K_a = 2$  between 9 and 40 GHz have now been measured. Spectroscopic constants were determined by fitting a theoretical spectrum calculated from a standard asymmetric-top Hamiltonian with centrifugal distortion to the observed frequencies. With three rotational constants and all five quartic centrifugal distortion constants (see Table 1), a fit with an rms of only 1 kHz is achieved, in excellent agreement with the 2–4 kHz measurement uncertainty. Sixteen lines of  ${}^{34}SS_2$  and nine lines of  ${}^{34}SS$  have also been detected (the larger number of lines for the former because  ${}^{34}S$  substitution at one of the equivalent sulfur atoms breaks the  $C_{2\nu}$  symmetry, enabling transitions normally symmetry-forbidden to be

Table 1. Rotational Constants of Isotopic S<sub>3</sub> (in MHz)<sup>a</sup>

constant	SSS	S <sup>34</sup> SS	<sup>34</sup> SSS
Α	23972.581(1)	23033.906(2)	23738.354(2)
В	2948.5467(3)	2948.6856(6)	2861.0127(4)
С	2622.2913(2)	2610.6946(5)	2550.1327(4)

<sup>*a*</sup> 1 $\sigma$  uncertainties (in parentheses) are in the units of the last significant digits. The best-fit centrifugal distortion constants (in MHz) for the normal isotopic species are:  $D_{\rm J} = 0.797(3) \times 10^{-3}, D_{\rm JK} = -20.33(5) \times 10^{-3}, D_{\rm K}$  $= 0.5033(2), d_1 = -0.1628(3) \times 10^{-3}, and d_2 = -6(1) \times 10^{-6}.$ 

observed). A more complete account of the laboratory measurements and data analysis will appear elsewhere.

Because of the symmetry of S<sub>3</sub>, only two geometrical parameters are needed to characterize its structure. An empirical  $(r_0)$  structure was derived (Figure 1) by adjusting the length of the S-S bond and the apex angle to reproduce the three measured rotational constants of the three isotopic species. The length of the S-S bond (1.917 Å) is close to that of  $S_2$  (1.889 Å), indicating substantial double-bonding. The angle (117.4°) is within 2° of that for the isovalent molecules O<sub>3</sub> (116.8°), SO<sub>2</sub> (119.5°), and S<sub>2</sub>O (118°), indicating sp<sup>2</sup> hybridization for the apex sulfur atom.<sup>17</sup> The separation between the terminal S atoms (3.150 Å) agrees fairly well with the van der Waals distance between S atoms in the solid allotropes, confirming that S<sub>3</sub> is best described as a bent chain rather than a three-membered ring. A structure incorporating theoretical corrections for zero-point vibration will be presented elsewhere.12

Starting with sulfur vapor in neon, in excess of  $5 \times 10^{12} \text{ S}_3$ clusters are produced in our molecular beam per gas pulse. This abundance, derived from line intensities relative to those of the stable molecule OCS at a known fractional abundance (i.e., 1% OCS in Ne), is comparable to that achieved for Si<sub>3</sub>  $(10^{13}/\text{pulse})$ .<sup>14</sup> At this high concentration, the high-resolution infrared and optical spectra of S<sub>3</sub> should be detectable with fairly standard spectroscopic techniques. The low-resolution spectrum near 410 nm that has been attributed to the  $C_{2\nu}$  form<sup>8</sup> may serve as a starting point for optical experiments with our rotational cold ( $T_{\rm rot} \approx 3$  K) molecular beam.

Thiozone may be abundant on Io, where S<sub>2</sub> and SO<sub>2</sub> are readily observed,<sup>18</sup> and it may be detectable in the interstellar gas as well. Because of widespread volcanism, S<sub>2</sub> and SO<sub>2</sub> are ejected in large quantities into the atmosphere of Io, where they drive a rich sulfur chemistry that has long fascinated planetary scientists. The red deposit from the plume of the volcano Pele is claimed to come from ultraviolet photolysis and polymerization of S2 to S3 and S4.2.18 S<sub>3</sub> is also of interstellar interest: about 10% of the 130 molecules detected in space contain sulfur; SiS, a molecule with two secondrow elements, is so abundant in the circumstellar shell of the carbon star IRC+10216 that six rare isotopic species have been detected.<sup>19</sup> With the precise rest frequencies that the present work provides, dedicated searches with large radio telescopes can now be undertaken for S<sub>3</sub> in the solar system and throughout the Galaxy. A list of the most relevant astronomical lines will be given elsewhere.

Because of the high concentration achieved here for S<sub>3</sub>, larger sulfur clusters may be detectable with the present technique. Ab initio calculations predict that S4 and S5 have singlet ground or extremely low-lying electronic states that are polar, with dipole moments comparable to or larger than that of S<sub>3</sub>. The low-lying polar state of S<sub>4</sub> is calculated to have a dipole moment of 0.87 D,<sup>20</sup> a planar *cis* structure with  $C_{2v}$  symmetry,<sup>1,10</sup> and a number of accessible rotational transitions.

Acknowledgment. We thank J. F. Stanton, R. J. McMahon, J. K. G. Watson, R. Steudel, L. S. Andrews, M. E. Jacox, H. S. P. Müller, and W. Klemperer for helpful discussions, and E. S. Palmer for laboratory assistance. This work is supported by NASA grant NAG5-9379 and NSF grant AST-9820722. S.T. is grateful to the Alexander von Humboldt Foundation for a Feodor Lynen research fellowship.

Note Added in Proof. The millimeter- and submillimeter-wave spectra of S<sub>3</sub> have now been detected, confirming the present geometrical structure, and yielding a better set of centrifugal distortion constants. The key rotational frequencies of S<sub>3</sub> into the far-IR can now be calculated to better than 1 ppm, more than adequate for an astronomical search. Following completion of the present work, S<sub>4</sub> was detected here by FTM spectroscopy, and its geometrical structure was determined to high precision by isotopic substitution.

## References

- (1) Ravaghavachari, K. J. Chem. Phys. 1990, 93, 5862-5874
- (2) (a) Moses, J. I.; Zolotov, M. Y.; Fegley, B., Jr. Icarus 2002, 156, 76-106. (b) Lellouch, E. Icarus 1996, 124, 1-24 and references therein.
- (a) Lewis, J. S.; Kreimendahl, F. A. *Icarus* 1980, 42, 330–337. (b) Prinn, R. G. *Geophys. Res. Lett.* 1979, 6, 807–810.
   (4) Glassgold, A. E.; Mamon, G. A. In *Chemistry and Spectroscopy of Interstellar Molecules*; Bohme, D. K., Herbst, E., Kaifu, N., Saito, S., Eds.; University Toyko Press: Toyko, 1992; pp 261–265.
- (5) Trambarulo, R.; Ghosh, S. N.; Burrus C. A., Jr.; Gordy, W. J. Chem. Phys. 1953, 21, 851-855
- (6) Eckert, B.; Steudel, R. Top. Curr. Chem. 2003, 231, 31-98.
- (a) Hassanzadeh, P.; Andrews, L. J. Phys. Chem. 1992, 96, 6579-6585. (b) Brabson, G. D.; Mielke, Z.; Andrews, L. J. Phys. Chem. 1991, 95, 79 - 86
- (8) Meyer, B.; Stroyer-Hansen, T.; Oommen, T. V. J. Mol. Spectrosc. 1972, 42, 335-343.
- (9) The S-S bond length was estimated to be 1.90  $\pm$  0.05 Å from photoelectron measurements (Nimlos, N. M.; Ellison, G. B. J. Phys. Chem. **1986**, 90, 2574–2580), and the apex angle was estimated to be  $116 \pm 2^{\circ}$ from the infrared matrix spectroscopy described in ref 7b.
- (a) Jones, R. O.; Ballone, P. J. Chem. Phys. 2003, 118, 9257-9265. (b) (10)Koch, W.; Natterer, J.; Heinemann, C. J. Chem. Phys. 1995, 102, 6159-6167
- (11)(a) Rice, J. E.; Amos, R. D.; Handy, N. C.; Lee, T. J.; Schaefer, H. F., III. J. Chem. Phys. 1986, 85, 963–968. (b) Millefioir, S.; Alparone, A. J. Phys. Chem. A 2001, 105, 9489-9497.
- (12) Stanton, J. F., personal communication, 2003. The dipole moment and cubic force field have been calculated at the CCSD(T)/cc-pVTZ level of theory.
- (13) Murray, C. W.; Handy, N. C.; Amos, R. D. J. Chem. Phys. 1993, 98, 7145-7151 and references therein.
- (14) McCarthy, M. C.; Thaddeus, P. Phys. Rev. Lett. 2003, 90, 213003/1-213003/4
- (15) Thaddeus, P.; McCarthy, M. C.; Thorwirth, S. Proceedings of the 4th Cologne-Bonn-Zermatt-Symposium "The Dense Interstellar Medium in
- (a) Yamamoto, S.; Saito, S. *Can. J. Phys.* **1994**, 72, 954–962. (b) Winnewisser, G.; Winnewisser, M.; Gordy, W. *J. Chem. Phys.* **1968**, 49, 0465 (2007). (16)3465-3478. (c) Behnke, M. Dissertation, Universität zu Köln, 2001.
- (17) Inorganic Chemistry; Wiberg, N., Holleman, A., Wilberg, E., Eds.; Academic Press: San Diego, 2001; p 532.
  (18) Spencer, J. R.; Jessup, K. L.; McGrath, M. A.; Ballester, G. E.; Yelle, R. Science 2000, 288, 1208–1210.
- (19) Cernicharo, J.; Guélin, M.; Kahane, C. Astron. Astrophys. Suppl. 2000, 142, 181-215
- (20) Quelch, G. E.; Schaefer, H. F.; Marsden, C. J. Am. Chem. Soc. 1990, 112, 8719-8733.

JA049645F